

Integrals Involving Bessel Functions

By

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Abstract

In this paper, I have evaluated some double integrals involving product of five Bessel functions of different arguments in terms of Appell's function F_4 with the help of operational calculus. The results obtained are believed to be new.

1. The following formulas are to be established here

$$\begin{aligned} & \int_0^\infty \int_0^\infty yx^{\mu_1+1} J_{\mu_1}(ax) K_{\mu_2}(b\sqrt{x^2+y^2}) I_{\mu_2}(c\sqrt{x^2+y^2}) J_{\mu_3}(gy) J_{\mu_3}(hy) dx dy \\ &= \frac{(a)^{\mu_1} (bc)^{\mu_2} (gh)^{\mu_3} 2^{\mu_1+1} \Gamma(\mu_1+\mu_2+\mu_3+2)}{(a^2+b^2+c^2+g^2+h^2)^{\mu_1+\mu_2+\mu_3+2}} \times \\ & F_4 \left[\frac{1}{2}(\mu_1+\mu_2+\mu_3+2), \frac{1}{2}(\mu_1+\mu_2+\mu_3+3); 1+\mu_2, 1+\mu_3; \right. \\ & \left. \frac{4b^2c^2}{(a^2+b^2+c^2+g^2+h^2)^2}, \frac{4g^2h^2}{(a^2+b^2+c^2+g^2+h^2)^2} \right], \end{aligned}$$

valid for $R(\mu_1+\mu_2+\mu_3+2) > 0$, $R(b) > |R(c)| + |I_m(g)| + |I_m(h)|$,

$$R(b) > |R(c)| + |I_m(g)|. \quad \dots\dots(1)$$

$$\begin{aligned} & \int_0^\infty \int_0^\infty yx^{\mu_1+1} J_{\mu_1}(ax) K_{\mu_2}(b\sqrt{x^2+y^2}) I_{\mu_2}(c\sqrt{x^2+y^2}) I_{\mu_3}(gy) I_{\mu_3}(hy) dx dy \\ &= \frac{(a)^{\mu_1} (bc)^{\mu_2} (gh)^{\mu_3} 2^{\mu_1+1} \Gamma(\mu_1+\mu_2+\mu_3+2)}{(a^2+b^2+c^2-g^2-h^2)^{\mu_1+\mu_2+\mu_3+2}} \times \\ & F_4 \left[\frac{1}{2}(\mu_1+\mu_2+\mu_3+2), \frac{1}{2}(\mu_1+\mu_2+\mu_3+3); 1+\mu_2, 1+\mu_3; \right. \\ & \left. \frac{4b^2c^2}{(a^2+b^2+c^2-g^2-h^2)^2}, \frac{4g^2h^2}{(a^2+b^2+c^2-g^2-h^2)^2} \right], \end{aligned}$$

valid for

$$R(\mu_1 + \mu_2 + \mu_3 + 2) > 0, \quad R(b) > |R(c)| + |R(g)| + |R(h)|,$$

$$R(b) > R(c) + |I_m(a)|. \quad \dots\dots(2)$$

$$\int_0^\infty \int_0^\infty x^{\mu_1+1} y^{2\mu_4-1} J_{\mu_1}(ax) K_{\mu_2}(b\sqrt{x^2+y^2}) I_{\mu_2}(c\sqrt{x^2+y^2}) \times$$

$$S_2 \left[\frac{1}{2}(\mu_3-1), \quad -\frac{1}{2}(\mu_3+1), \quad \frac{1}{2}(1-\mu_4), \quad -\frac{1}{2}\mu_4; \quad \frac{gy^2}{4} \right] dx dy$$

$$= \sum_{\mu_3, -\mu_3} \frac{(g)^{\mu_3} (bc)^{\mu_2} (a)^{\mu_1} \Gamma(-\mu_3) \Gamma(\mu_1 + \mu_2 + \mu_3 + \mu_4) 2^{\mu_1 + \mu_3 + 3\mu_4 - 4}}{\sqrt{\pi} \Gamma(1 + \mu_2) (a^2 + b^2 + c^2)^{\mu_1 + \mu_2 + \mu_3 + \mu_4}} \times$$

$$F_4 \left[\frac{1}{2}(\mu_1 + \mu_2 + \mu_3 + \mu_4), \quad \frac{1}{2}(\mu_1 + \mu_2 + \mu_3 + \mu_4 + 1); \quad 1 + \mu_2, \quad 1 + \mu_3; \right.$$

$$\left. \frac{4b^2c^2}{(a^2 + b^2 + c^2)^2}, \quad \frac{16g^2}{(a^2 + b^2 + c^2)^2} \right],$$

valid for

$$R(\mu_1 + \mu_2 + \mu_3 + \mu_4) > 0, \quad R(b) > |R(c)|, \quad R(b) > |R(c)| + |I_m(a)|, \quad g > 0. \dots(3)$$

$$\int_0^\infty \int_0^\infty x^{\mu_1+1} y^{2\mu_3+2\mu_4+3} J_{\mu_1}(ax) K_{\mu_2}(b\sqrt{x^2+y^2}) I_{\mu_2}(c\sqrt{x^2+y^2})$$

$$\times {}_1F_2 \left(\frac{1}{2} + \mu_3; \quad 2\mu_3 + 1, \quad \mu_3 + \mu_4; \quad -g^2 y^2 \right) dx dy$$

$$= \frac{(a)^{\mu_1} (bc)^{\mu_2} \Gamma(\mu_3 + \mu_4) \Gamma(\mu_1 + \mu_2 + \mu_3 + \mu_4 + 1) 2^{2\mu_2 + 2\mu_3 + 2\mu_4 - 5}}{\Gamma(1 + \mu_2) (a^2 + b^2 + c^2 + g^2)^{\mu_1 + \mu_2 + \mu_3 + \mu_4 + 1}} \times$$

$$F_4 \left[\frac{1}{2}(\mu_1 + \mu_2 + \mu_3 + \mu_4 + 1), \quad \frac{1}{2}(\mu_1 + \mu_2 + \mu_3 + \mu_4 + 2); \right.$$

$$\left. 1 + \mu_2, \quad 1 + \mu_3; \quad \frac{4b^2c^2}{(a^2 + b^2 + c^2 + g^2)^2}, \quad \frac{4g^2}{(a^2 + b^2 + c^2 + g^2)^2} \right],$$

valid for

$$R(\mu_3 + \mu_4) > 0, \quad R(\mu_1 + \mu_2 + \mu_3 + \mu_4 + 1) > 0, \quad R(b) > |R(c)| + |I_m(g)|,$$

$$R(b) > |R(c)| + |I_m(a)|. \quad \dots\dots(4)$$

For the definition of Appells' function F_4 , see [1, p. 224].

2. As usual $\Phi(p) \doteq h(t)$ will be used to denote Laplace transform defined by the equation

$$\Phi(p) = p \int_0^\infty e^{-pt} h(t) dt, \quad \dots\dots(5)$$

where the integral is convergent and $R(p) > 0$.

We prove a theorem of Laplace transform, which will be used in our investigation.

THEOREM: *If $h(t) \in L(0, \infty)$ and if $R(\mu_1) > -1$, $a > 0$, $R(p) > 0$, then*

$$t^{-\mu_1-1} e^{-\frac{a}{t}} h(t) \doteq p a^{-\frac{1}{2}\mu_1} \int_0^\infty x^{\frac{1}{2}\mu_1} (x+p)^{-1} J_{\mu_1}(2\sqrt{ax}) \Phi(x+p) dx,$$

where $h(t) \doteq \Phi(p)$(6)

PROOF: We know that if

$$h(t) \doteq \Phi(p)$$

then, Erdélyi [1, p. 129],

$$e^{-\alpha t} h(t) \doteq p(p+\alpha)^{-1} \Phi(p+\alpha). \quad \dots\dots(7)$$

Also, Erdélyi [1, p. 245],

$$t^{\frac{1}{2}\mu_1} J_{\mu_1}(2\sqrt{at}) \doteq p^{-\mu_1} a^{\frac{1}{2}\mu_1} e^{-a/p}. \quad \dots\dots(8)$$

Using the relations (7) and (8) in Goldstiens' result [2] that if

$$h_1(p) \doteq g_1(t) \text{ and } h_2(p) \doteq g_2(t),$$

then $\int_0^\infty h_1(t) g_2(t) t^{-1} dt = \int_0^\infty h_2(t) g_1(t) t^{-1} dt$,(9)

and replacing α by p , we have

$$\int_0^\infty t^{-\mu_1-1} e^{-pt-\frac{a}{t}} h(t) dt = a^{-\frac{1}{2}\mu_1} \int_0^\infty x^{\frac{1}{2}\mu_1} (x+p)^{-1} J_{\mu_1}(2\sqrt{ax}) \Phi(x+p) dx. \quad \dots(10)$$

The theorem follows immediately on interpreting the left side of (10).

To prove (1), we take, Erdélyi [1, p. 284],

$$\begin{aligned} h(t) &= t^{-1} \exp\left[-\frac{b+c}{t}\right] I_{\mu_2}\left(\frac{2\sqrt{bc}}{t}\right) \\ &\doteq 2p K_{\mu_2}(2\sqrt{bp}) I_{\mu_2}(2\sqrt{cp}) \end{aligned}$$

$$= \emptyset(p), R(p) > 0, R(b) > 0, R(c) > 0. \quad \dots\dots(11)$$

Using (11) in (6), we have

$$\begin{aligned} & t^{-\mu_1-2} \exp\left[-\frac{a+b+c}{t}\right] I_{\mu_2}\left(\frac{2\sqrt{bc}}{t}\right) \\ & \doteq 2p a^{-\frac{1}{2}\mu_1} \int_0^\infty x^{\frac{1}{2}\mu_1} J_{\mu_1}(2\sqrt{ax}) K_{\mu_2}[2\sqrt{b(x+p)}] I_{\mu_2}[2\sqrt{c(x+p)}] dx. \quad \dots\dots(12) \end{aligned}$$

Also, Erdélyi [1, p. 281],

$$J_{\mu_3}(2\sqrt{gt}) J_{\mu_3}(2\sqrt{ht}) \doteq \exp\left[-\frac{g+h}{t}\right] I_{\mu_3}\left(\frac{2\sqrt{gh}}{t}\right), \quad \dots\dots(13)$$

where $R(\mu_3) > -1$.

Applying (12) and (13) in (9), we have

$$\begin{aligned} & \int_0^\infty \int_0^\infty x^{\frac{1}{2}\mu_1} J_{\mu_1}(2\sqrt{ax}) K_{\mu_2}[2\sqrt{b(x+t)}] I_{\mu_2}[2\sqrt{c(x+t)}] J_{\mu_3}(2\sqrt{gt}) J_{\mu_3}(2\sqrt{ht}) dt dx \\ & = \frac{a^{\frac{1}{2}\mu_1}}{2} \int_0^\infty t^{\mu_1+1} \exp[-(a+b+c+g+h)t] I_{\mu_2}(2\sqrt{bc}t) I_{\mu_3}(2\sqrt{gh}t) dt. \end{aligned}$$

(1) follows immediately on evaluating the single integral with the help of the result, Erdélyi [1, p. 196].

(2), (3), (4) can be proved in the same way by using the results [1, p. 281, 282, 284],

$$I_\rho(2a\sqrt{t}) I_\rho(2b\sqrt{t}) \doteq p^{-1} e^{(a^2+b^2)/p} I_\rho\left(\frac{2ab}{p}\right), \quad \dots\dots(14)$$

where $R(\rho) > -1$ and $R(p) > 0$,

$$\begin{aligned} & t^{-2\lambda-1} S_2\left(\rho - \frac{1}{2}, -\rho - \frac{1}{2}, \lambda + \frac{1}{2}, \lambda; \frac{1}{2}at\right) \\ & \doteq \frac{p^{2\lambda}}{2^{2\lambda}\sqrt{\pi}} K_{2\rho}\left(\frac{2a}{p}\right), \quad \dots\dots(15) \end{aligned}$$

where $R(\lambda \pm \rho) < 0$, and

$$\begin{aligned} & t^{\lambda+v-1} {}_1F_2\left(v + \frac{1}{2}; 2v+1, \lambda+v; -2at\right) \\ & \doteq \frac{\Gamma(v+1)\Gamma(\lambda+v)}{2^{-v} a^v p^\lambda} e^{-\frac{a}{p}} I_v\left(\frac{a}{p}\right), \quad \dots\dots(16) \end{aligned}$$

where $R(\lambda + \nu) > 0$, instead of (13) respectively.

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References

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- [2] Goldstien, S., Operational representation of Whittakers' confluent hypergeometric function and Weber's parabolic cylinder functions, Proc. Lond. Math. Soc. **34**, 103-125, (1932).