On the Remainder Terms of the Optimal Formulas

By

Yoshitane Shinohara (Received September 30, 1967)

Abstract

The remainder term of the optimal multi-step formula, that is the stable formula of maximum order, is considered from the analytical point of view.

§1. Introduction

For the numerical integration of ordinary differential equation

$$\frac{dy}{dx} = f(x, y),$$

we consider the multi-step formula

(1.2)
$$\sum_{j=0}^{n} \alpha_{j} y(x_{k+j}) = h \sum_{j=0}^{n} \beta_{j} f(x_{k+j}, y(x_{k+j})) + Ch^{p+1} y^{(p+1)}(\xi) \quad (x_{k} \leq \xi \leq x_{k+n}),$$

where y = y(x) is a solution of (1.1) on the interval [a, b] satisfying the initial condition $y(a) = y_0$ and we shall always assume that $\alpha_n = 1$, $|\alpha_0| + |\beta_0| > 0$. If we associate with the formula (1.2) the polynomials

(1.3)
$$\begin{cases} \rho(\zeta) = \alpha_n \zeta^n + \alpha_{n-1} \zeta^{n-1} + \dots + \alpha_1 \zeta + \alpha_0, \\ \sigma(\zeta) = \beta_n \zeta^n + \beta_{n-1} \zeta^{n-1} + \dots + \beta_1 \zeta + \beta_0, \end{cases}$$

where the polynomials $\rho(\zeta)$ and $\sigma(\zeta)$ have no common factor, then, the stability conditions of (1.2) can be written as follows:

$$(1.4) \qquad \qquad \text{if } \rho(\zeta)\!=\!0, \text{ then } \begin{cases} \text{either} & |\zeta|\!<\!1\\ \text{or} & |\zeta|\!=\!1 \text{ and } \rho'(\zeta)\!\neq\!0. \end{cases}$$

Moreover, if we postulate the conditions

$$\alpha_{\nu} = -\alpha_{n-\nu}, \quad \beta_{\nu} = \beta_{n-\nu},$$

then, the order p of the local truncation error of the stable multi-step method (1.2) becomes p=n+2 for even n, by the theorem of Dahlquist [1]. Hence, from (1.2), (1.3), (1.4) and (1.5), we have

where $y_n = y(x_n)$, $f_n = f(x_n, y(x_n))$ and E is the operator such that Ey(x) = y(x+h). The stable multi-step formula (1.6) of maximum order will be called an optimal formula $\lceil 3 \rceil$.

In the present paper, first, we express the constant C explicitly in terms of the coefficients α_i . Next, we have obtained the optimal formulas for n=4 and n=6, and then we show that the optimal formula for which the quantity |C| attains a minimum does not exist for n=4 and n=6. Lastly, we compare them with Henrici's optimal formula [3] and Gorbunov's optimal formula [2] in a numerical example.

$\S 2$. Determination of the constant C

As is well-known, we have $E=e^{hD}$ formally where D=d/dx. Hence, from (1.6), we have

$$\rho(e^{hD}) - h\sigma(e^{hD})D \sim Ch^{n+3}D^{n+3} \quad (h \rightarrow 0)$$

formally. Replacing the operator e^{hD} by the scalar ζ , we see that the above relation is equivalent to the following one:

(2.1)
$$\rho(\zeta) - (\log \zeta)\sigma(\zeta) \sim C(\zeta - 1)^{n+3} \quad (\zeta \to 1),$$

where $\log \zeta$ is a branch such that $\log \zeta \rightarrow 0$ as $\zeta \rightarrow 1$.

If we put $\eta = \zeta - 1$, then (2.1) can be rewritten as follows:

(2.2)
$$\frac{\rho(1+\eta)}{\log(1+\eta)} - \sigma(1+\eta) \sim C\eta^{n+2} \quad (\eta \to 0).$$

According to the conditions (1.5) of maximum order, we may write:

$$\begin{split} \rho(\zeta) &= (\zeta^{n} - 1)\alpha_{0} + (\zeta^{n-2} - 1)\zeta\alpha_{1} + \dots + (\zeta^{2} - 1)\zeta^{\frac{n}{2} - 1}\alpha_{\frac{n}{2} - 1} \\ &= \sum_{i=0}^{\frac{n}{2} - 1} \zeta^{i}(\zeta^{n-2i} - 1)\alpha_{i} \\ &= \sum_{i=0}^{\frac{n}{2} - 1} \zeta^{i}(\zeta^{2(\frac{n}{2} - i)} - 1)\alpha_{i} \end{split}$$

$$=\sum_{i=0}^{\frac{n}{2}-1} \zeta^i (\zeta^2-1) (\zeta^{n-2i-2}+\zeta^{n-2i-4}+\dots+\zeta^2+1) \alpha_i.$$

Hence, we have

(2.3)
$$\rho(1+\eta) = \sum_{i=0}^{\frac{n}{2}-1} \eta(2+\eta)(1+\eta)^{i} \left[1 + (1+\eta)^{2} + \dots + (1+\eta)^{n-2(i+1)}\right] \alpha_{i}$$

$$= \sum_{i=0}^{\frac{n}{2}-1} \sum_{j=0}^{\frac{n}{2}-1-i} \eta(2+\eta)(1+\eta)^{i} (1+\eta)^{2j} \alpha_{i}$$

$$= \sum_{i=0}^{\frac{n}{2}-1} \sum_{j=0}^{\frac{n}{2}-1-i} \eta(2+\eta)(1+\eta)^{i+2j} \alpha_{i}.$$

Then, since

$$\frac{\eta(1+\eta)^{r-1}}{\log(1+\eta)} = \int_0^1 (1+\eta)^{x+r-1} dx,$$

we have, from (2.3),

(2.4)
$$\frac{\rho(1+\eta)}{\log(1+\eta)} = (2+\eta) \sum_{i=0}^{\frac{n}{2}-1} \alpha_i \sum_{j=0}^{\frac{n}{2}-1-i} \int_0^1 (1+\eta)^{x+i+2j} dx.$$

Since $\sigma(1+\eta)$ is a polynomial of degree n at most, the constant C must be the coefficient of η^{n+2} in the power series $\rho(1+\eta)/\log(1+\eta)$ at $\eta=0$. Therefore, we obtain from (2.2) and (2.4),

(2.5)
$$C = \sum_{i=0}^{\frac{n}{2}-1} \alpha_i \sum_{j=0}^{\frac{n}{2}-1-i} \left\{ \frac{2}{(n+2)!} \int_0^1 (x+i+2j)^{\lfloor n+2 \rfloor} dx + \frac{1}{(n+1)!} \int_0^1 (x+i+2j)^{\lfloor n+1 \rfloor} dx \right\}.$$

A simple calculation gives for instance the following values of the constant C for n=2, 4, 6:

(2.6)
$$\begin{cases} n = 2, & C = \frac{1}{90} \alpha_0 \\ n = 4, & C = \frac{1}{3780} (5\alpha_1 + 32\alpha_0) \\ n = 6, & C = \frac{1}{907200} (832\alpha_1 - 184\alpha_2 + 5832\alpha_0) \end{cases}$$

where $\alpha_0 = -1$.

§3. Table of the optimal formulas and a numerical example

Taking into account the properties $\alpha_n=1$, $\alpha_{\nu}=-\alpha_{n-\nu}$, $\beta_{\nu}=\beta_{n-\nu}$, (2.6) and stability condition (1.4), we have

n=2:

$$y_{k+2} = y_k + \frac{1}{3}h(f_{k+2} + 4f_{k+1} + f_k) - \frac{1}{90}h^5y^{(5)}(\xi)$$

(Simpson's rule),

n=4:

(3.1)
$$y_{k+4} + \mu (y_{k+3} - y_{k+1}) - y_k$$

$$= h(\gamma_1 f_{k+4} + \gamma_2 f_{k+3} + \gamma_3 f_{k+2} + \gamma_2 f_{k+1} + \gamma_1 f_k)$$

$$+ \frac{5\mu - 32}{3780} h^7 y^{(7)}(\xi),$$

where

$$\mu = \alpha_1,$$
 $\gamma_1 = \frac{1}{1+\lambda^2} \left(\frac{1}{3} \lambda^2 + \frac{13}{45} \right),$
 $\gamma^2 = \frac{1}{1+\lambda^2} \left(\frac{2}{3} \lambda^2 + \frac{98}{45} \right),$
 $\gamma_3 = \frac{1}{1+\lambda^2} \left(-2\lambda^2 + \frac{138}{45} \right),$
 $\lambda^2 = \frac{2-\mu}{2+\mu} \qquad (|\mu| < 2).$

For n = 6, we have

$$(3.2) y_{k+6} + \mu (y_{k+5} - y_{k+1}) + \lambda (y_{k+4} - y_{k+2}) - y_k$$

$$= h(\delta_1 f_{k+6} + \delta_2 f_{k+5} + \delta_3 f_{k+4} + \delta_4 f_{k+3} + \delta_3 f_{k+2} + \delta_2 f_{k+1} + \delta_1 f_k)$$

$$- \frac{1}{907200} (184\lambda - 832\mu + 5832)h^9 y^{(9)}(\xi),$$

where

$$\begin{split} \mu &= \alpha_1, \quad \lambda = \alpha_2, \\ \delta_1 &= \frac{1}{1 + \alpha^2 + \beta^2} \left\{ \frac{1}{2} + \frac{1}{6} (3\alpha^2 - 1) + \frac{1}{90} (45\beta^2 - 15\alpha^2 - 4) \right. \\ &\quad \left. - \frac{1}{1890} (84\alpha^2 + 315\beta^2 + 44) \right\}, \\ \delta_2 &= \frac{1}{1 + \alpha^2 + \beta^2} \left\{ 3 + \frac{1}{3} (3\alpha^2 - 1) - \frac{1}{45} (45\beta^2 - 15\alpha^2 - 4) \right. \\ &\quad \left. + \frac{1}{315} (84\alpha^2 + 315\beta^2 + 44) \right\}, \\ \delta_3 &= \frac{1}{1 + \alpha^2 + \beta^2} \left\{ \frac{15}{2} - \frac{1}{6} (3\alpha^2 - 1) - \frac{1}{90} (45\beta^2 - 15\alpha^2 - 4) \right. \\ &\quad \left. - \frac{1}{126} (84\alpha^2 + 315\beta^2 + 44) \right\}, \\ \delta_4 &= \frac{1}{1 + \alpha^2 + \beta^2} \left\{ 10 - \frac{2}{3} (3\alpha^2 - 1) + \frac{2}{45} (45\beta^2 - 15\alpha^2 - 4) \right. \\ &\quad \left. + \frac{2}{189} (84\alpha^2 + 315\beta^2 + 44) \right\}, \\ \alpha^2 &= \frac{40 - 8\lambda}{8\mu + 4\lambda + 12} , \\ \beta^2 &= \frac{-8\mu + 4\lambda + 12}{8\mu + 4\lambda + 12} . \end{split}$$

In this case, the stability domain D is as follows:

$$D\colon egin{array}{c} \lambda < rac{\mu^2}{4} + 1 \ \lambda > -2\mu - 3 \ \lambda \geqq 2\mu - 3 \ |\mu| < 4. \end{array}$$

The above analysis shows that the minimum value of the constant C in absolute value can not be obtained under the stability conditions in both

cases for n=4 and n=6, because $\mu=2$ (for n=4) and $\mu=4$, $\lambda=5$ (for n=6) are not contained in stability domain, respectively. The formula (3.1) for $\mu=1.41477653$ is Gorbunov's formula [2]. The formulas (3.2) for $\mu=-1$, $\lambda=1$ and $\mu=1.416369190$, $\lambda=1.002253240$ are Henrici's formula [3] and Gorbunov's formula [2], respectively.

We illustrate the optimal formula for which the error is smaller than that of the formulas obtained by Gorbunov and Henrici.

The Cauchy problem:

$$\frac{dy}{dx} = y$$
, $y(0) = 1$

is integrated numerically with a fixed step-size h = 0.2.

The starting values are computed from the exact solution e^x and the following predictor formulas are used for (3.1) and (3.2), respectively.

$$y_{k+4} - y_{k+3} = \frac{1}{24} h(55f_{k+3} - 59f_{k+2} + 37f_{k+1} - 9f_k),$$

$$y_{k+6} - y_{k+5} = \frac{1}{1440} h(4277f_{k+5} - 7923f_{k+4} + 9982f_{k+3} - 7298f_{k+2} + 2877f_{k+1} - 475f_k).$$

			10.0	
	C	$y_k - e^{x_k} (x_k = 10.6)$	μ	
$-7.142858300 \times 10^{-3}$ $-6.594210939 \times 10^{-3}$ $-6.481483217 \times 10^{-3}$ $-6.150795676 \times 10^{-3}$ $-5.820108134 \times 10^{-3}$		3. 10×10^{-2}	0. 9999991252	
		2.52×10^{-2}	1. 4147765300	
		2.42×10^{-2}	1.4999986880	
		2.14×10^{-2}	1. 7499984690	
		1.99×10^{-2}	1. 9999982500	
	n=6			
С	$y_k - e^{x_k} (x_k = 11)$	λ	μ	
7. 548500882×10	2.92×10^{-3}	1.000000000	-1.000000000	
$5.332887379 \times 10^{-1}$	6.98×10^{-4}	1. 002253240	1. 416369190	
$4.595442883 \times 10^{-}$	4.72×10^{-4}	1. 788854382	2. 394427191	

The computation is carried out in the floating-point arithmetic with 37 bits mantissa and rounding is done by chopping.

As is well-known, the optimal multi-step method is weakly unstable. From the practical point of view, therefore, the corrector formula such that

the quantity |C| of the remainder term is minimum under the strong stability should be considered. For a study of these we refer to Urabe's formulas [4, 5].

Department of Applied Mathematics Faculty of Engineering Tokushima University

References

- [1] Dahlquist, G., Convergence and stability in the numerical integration of ordinary differential equations, Math. Scand., 4 (1956), 33-53.
- [2] Gorbunov, A. D. and Šebalina, O. P., Predicting-correcting methods with optimum correction formula, Compt. Methods and Programming (Compt. Center Moscow Univ. Collect Works 2) (Russian), Izdat. Moscow Univ., Moscow (1965), 275-280.
- [3] Henrici, P., Discrete variable methods in ordinary differential equations, Wiley, 1962.
- [4] Shinohara, Y., On Urabe's predictor-corrector method for the numerical solution of ordinary differential equations, Bulletin of Faculty of Engineering, Tokushima Univ., 3-1 (1966), 75-85.
- [5] Urabe, M., Yanagiwara, H. and Shinohara, Y., Periodic solution of van der Pol's equation with damping coefficient λ=2-10, J. Sci. Hiroshima Univ., Ser. A, 23 (1960), 325-366.